

# Chiral dynamics in hard processes

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- Chiral symmetry breaking - the phenomenon shaping our World
- New soft pion theorems in hard processes– how ChSB works for colour clusters
- Data: present & possibilities to discover new chiral phenomena @ PANDA

# QCD & chiral symmetry

$$S_{QCD} = \int d^4x \left\{ -\frac{1}{4} G_{\mu\nu}^a G^{a\mu\nu} + \sum_f \bar{\Psi}_f (i\partial - g A^a t^a - m_f) \psi_f \right\}$$

Contains all strong interaction phenomena

$$m_p = \text{const} \Lambda_{QCD}$$

$$m_{U^{235}} = \text{const} \Lambda_{QCD}$$

$$\sigma_{tot}(p + U^{235}) = \frac{1}{\Lambda_{QCD}^2} f\left(\frac{\Lambda_{QCD}}{E}\right)$$

Transmutation of dimension: scale is given not by quark masses but

$$\Lambda_{QCD} = M_{cut} e^{-\frac{8\pi^2}{bg^2(M_{cut})}}$$

$$b = \frac{11}{3} N_c - \frac{2}{3} N_f$$

Symmetries:

$$m_d \ll m_d \ll 0$$

$$U_L(N_f) \otimes U_R(N_f)$$

But:  $m_\rho = 770 \text{ MeV}$

$$m_N = 940 \text{ MeV}$$

$$m_{\Lambda_1} \ll 1300 \text{ MeV}$$

$$m_{N^*} \ll 1535 \text{ MeV}$$

# Chiral symmetry is broken spontaneously

Nambu, Nobel Prize 2008

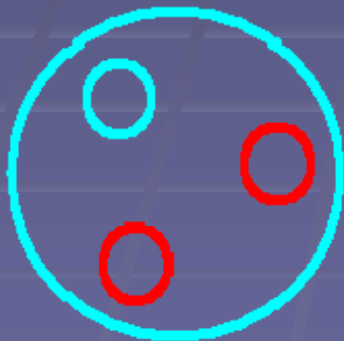
$\Rightarrow$  Goldstone bosons  $\leftarrow$  pions

Order parameters  $\leftarrow \langle 0 | \bar{\psi} \psi | 0 \rangle \simeq -(250 \text{ MeV})^3$

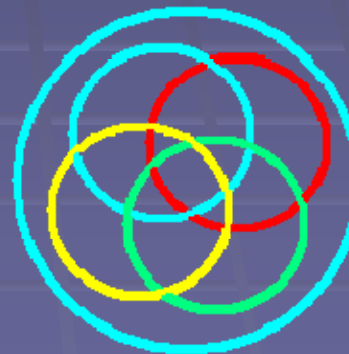
$\leftarrow$  dynamical quark mass

$M \simeq 350 \text{ MeV}$

One of puzzles: Why nucleon looks like



not like



Why  $MR_0 \ll 1$ ? Hidden small parameter?

To solve this and other puzzles one has to understand mechanisms of ChSB in QCD.  
As important as understanding of confinement!

Due to spontaneous chiral symmetry breakdown: 95% mass of the nucleon is generated → 95 % mass of visible Universe!

Chiral symmetry is broken spontaneously that is why:

- Protons and neutrons exist and are massive
- Pions are light
- Nuclei exist
- ... We are here and can study strong interactions @ FAIR

**Spontaneous breakdown of chiral symmetry really shapes our world!**

Up to now ChSB phenomenon has been studied only in soft processes and it was never studied at the level of quark and gluon degrees of freedom.

Main idea of soft pion theorems is in the following equation:

$$|\pi \text{ state}\rangle = i Q |0\rangle$$



generator of broken subgroup

$$\lim_{\kappa^\mu \rightarrow 0} \langle \pi^a(\kappa) B | \theta | A \rangle =$$

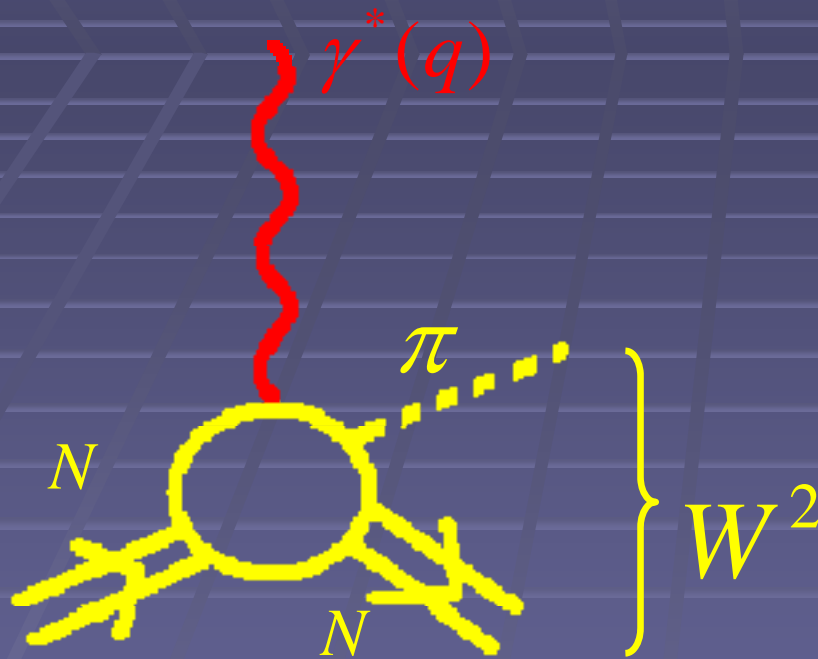
$$= \frac{i}{f_\pi} \langle B | [Q_5^a, \theta] | A \rangle + \dots$$

↑  
computable by symmetry  
transformation

*e.g.*

$$[Q_5^a, \psi(\kappa)] = \gamma_5 \tau^a \psi(x)$$

# Pion production near threshold



$$-q^2 = Q^2$$

$$W - W_{th} \ll m_\pi$$

$$W_{th} = M_N + m_\pi$$

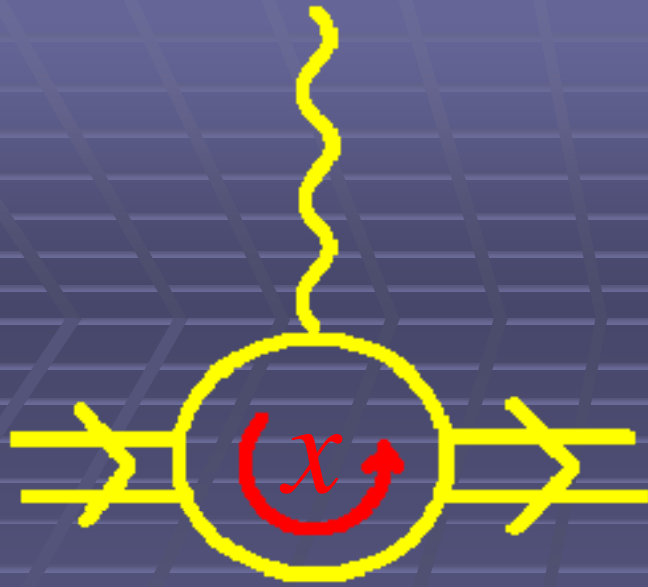
Pion is a (pseudo) Nambu-Goldstone Boson of  
SBChS  $\Rightarrow$  soft pion theorem

/Nambu, Lurie '62

Nambu, Shraumer '62/

These ideas created a new direction:

ChPT. Note however that the fundamental d.o.f. Of QCD are not involved!



- chirally rotated e.m. current  
 $\Rightarrow$  related to axial FF


Example:

$$A(\gamma^* p \rightarrow \pi^+ n) \Big|_{W=W_{th}} =$$

$$= \frac{1}{\sqrt{2}f_\pi} \left[ G_A(Q^2) + g_A G_{Mn}(Q^2) \frac{Q^2}{Q^2 + 2M_N} \right] + O\left(\frac{m_\pi}{\Lambda}\right)$$



$$\lim_{\kappa^\mu \rightarrow 0} \langle \pi^a(\kappa) B | \theta | A \rangle =$$

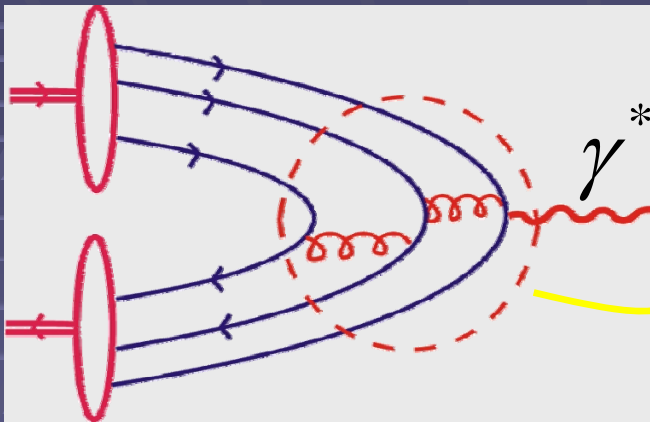
$$= \frac{i}{f_\pi} \langle B | [Q_5^a, \theta] | A \rangle + \dots$$


Would be breakthrough if here we would have a well defined cluster of quarks and gluons! Would allow to connect ChSB with fundamental d.o.f. of QCD! Big potential for discovery of new phenomena and understanding of mechanisms of “world shaping phenomenon” in QCD.

*e.g.*

$$[Q_5^a, \psi(\kappa)] = \gamma_5 \tau^a \psi(x)$$

In the limit of  $Q^2 \rightarrow \infty$  Nucleon FF's are described in terms of partons-clusters of quarks



hard interactions  
=>pQCD

Nucleon

Light-cone WF

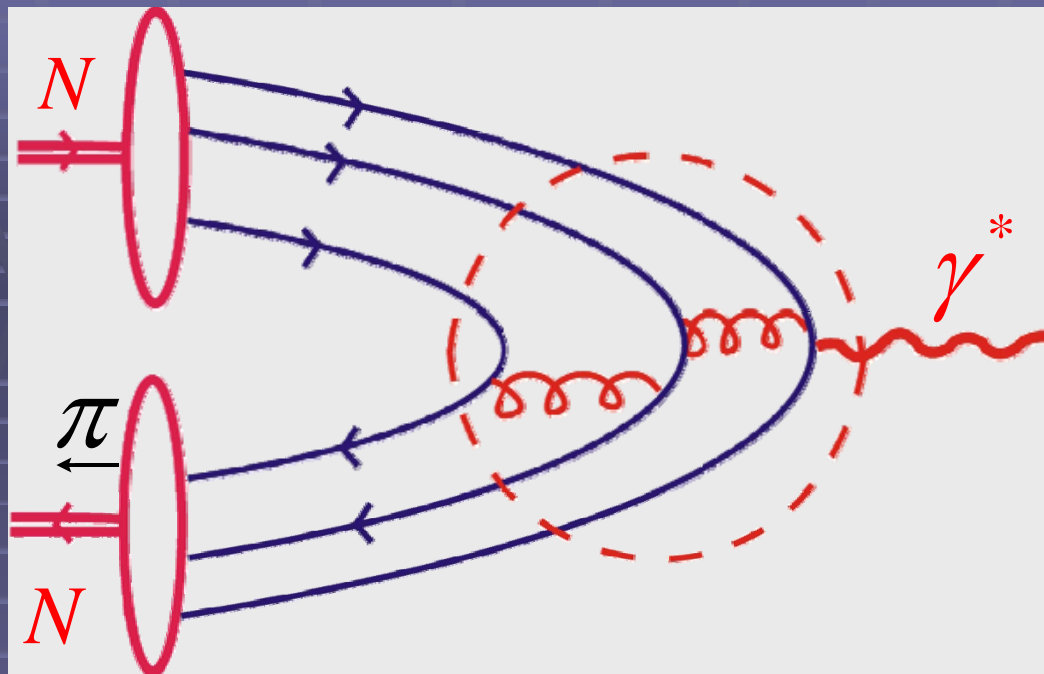
$$F(Q^2) \simeq \frac{\alpha_s^2(Q)}{Q^4} \int [dx][dy] \phi_N(x) T_H(x, y) \phi_N(y)$$

Filtering out the lowest Fock component of the hadron's WF

$$\alpha_s = \frac{g^2}{4\pi}$$

/S. Brodsky, P. Lepage/

/Efremov, Radyushkin/



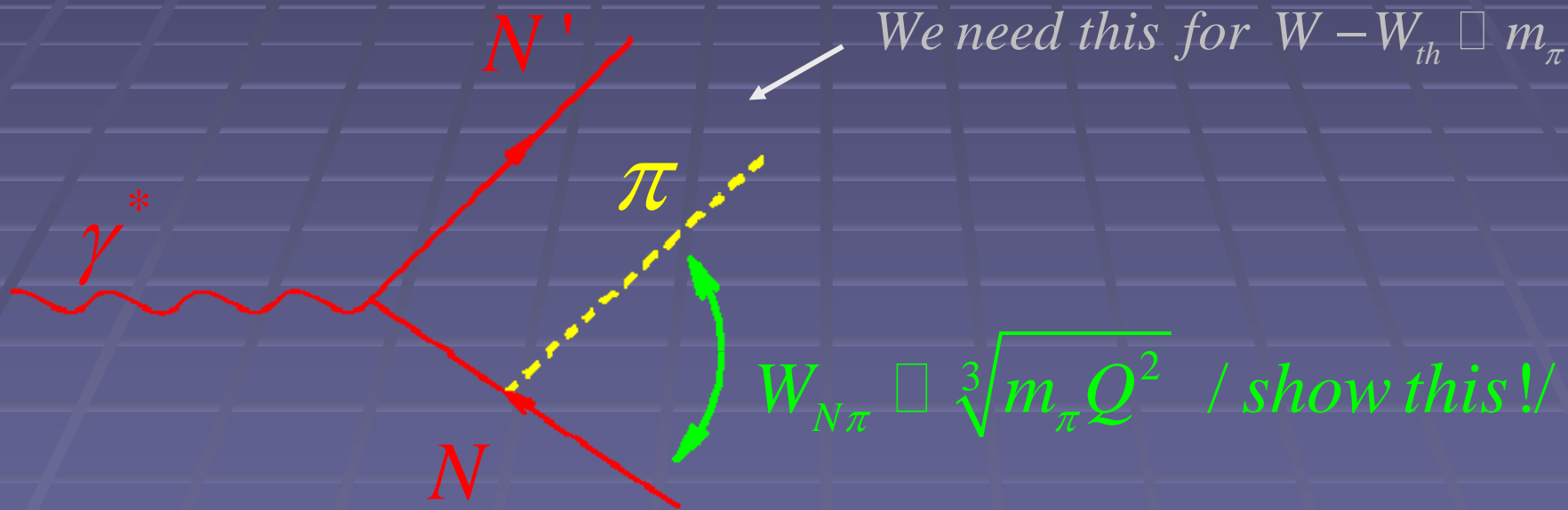
$$A(\gamma^* N \rightarrow \pi N)|_{th} \simeq \frac{\alpha_s^2}{Q^4} \int [dx][dy] \Phi_N(x) T_H(x, y) \Phi_{\pi N}(y)$$

LC WF of nucleon

LC WF of  $\pi N$  system

For NS theorem to work the pion should soft  
relative to all hadrons!

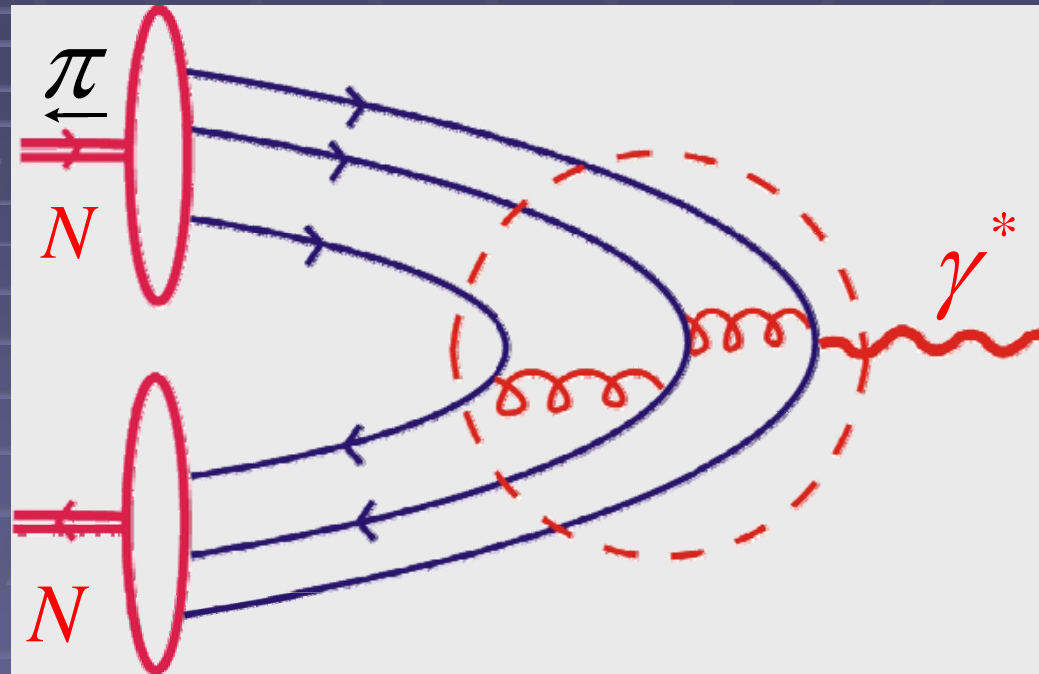
*Let us take first  $Q^2 \rightarrow \infty$  and then  $m_\pi \rightarrow 0$*



$$\Rightarrow \text{for } Q^2 \gg \frac{\Lambda^3}{m_\pi}$$

the pion is not soft relative to  
initial nucleon

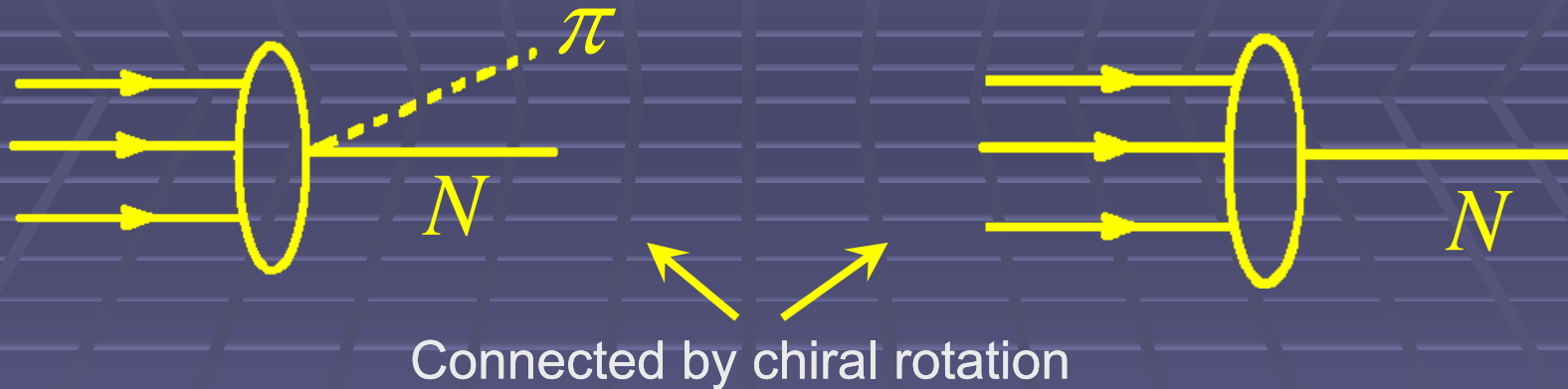
Higher twists ( $1/Q^2$  suppressed) – suppression of initial Nambu-Goldstone emission!



*For  $Q^2 \gg \frac{\Lambda^3}{m_\pi}$  one has prove SPT*

*For LC WF of  
 $\pi N$  Low mass system*

What are  $\pi N$  WF's when  $\pi$  is near threshold?



$$p \uparrow \rangle = \frac{\varphi_s}{\sqrt{6}} |2U \uparrow d \downarrow U \uparrow - U \uparrow U \downarrow d \uparrow - d \uparrow U \downarrow U \downarrow \rangle + \\ + \frac{\varphi_a}{\sqrt{2}} |U \uparrow U \downarrow d \uparrow - d \uparrow U \downarrow U \uparrow \rangle$$

$$p \uparrow \pi^0 \rangle = \frac{\varphi_s}{2\sqrt{6}f_\pi} |6U \uparrow d \downarrow U \uparrow + U \uparrow U \downarrow d \uparrow + d \uparrow U \downarrow U \uparrow \rangle + \\ + \frac{\varphi_a}{2\sqrt{2}f_\pi} |U \uparrow U \downarrow d \uparrow - d \uparrow U \downarrow U \uparrow \rangle$$

/Pobylitsa, Strikman, MVP, PRL 2001/

Simple algebra!

New theorem = “1/2” of NS theorem

For new SPTh important that QCD is  
a gauge theory!

## Comparison

"Nambu – Shrauner window"  $\left( \Lambda^2 \ll Q^2 \ll \frac{\Lambda^3}{m_\pi} \right)$

$$A(\gamma^* p \rightarrow \pi^+ n)|_{th} = \frac{1}{\sqrt{2} f \pi} (G_A + g_A G_{Mn})$$

"Our window"  $\left( Q^2 \gg \frac{\Lambda^3}{m_\pi} \gg \Lambda^2 \right)$

/Pobylitsa, Strikman, MVP, PRL 2001/

$$A(\gamma^* p \rightarrow \pi^+ n)|_{th} = \frac{1}{\sqrt{2} f \pi} \left( \frac{5}{3} G_{Mp} + \frac{4}{3} G_{Mn} \right)$$

What about data?



→ *E136 SLAC experiment*

$F_2^P(w, Q^2)$  at  $W - W_{th} \leq m_\pi$  &  $Q^2 \geq 9 \text{GeV}^2$

/P.Bosted et al. '94/

$F_2^P(W, Q^2) \underset{Q^2 \rightarrow \infty}{\sim} \frac{1}{Q^6}$  for both theorems

but the coefficient is given by different combinations of FF's of nucleon

# E136 experiment /Bosted et al '94/

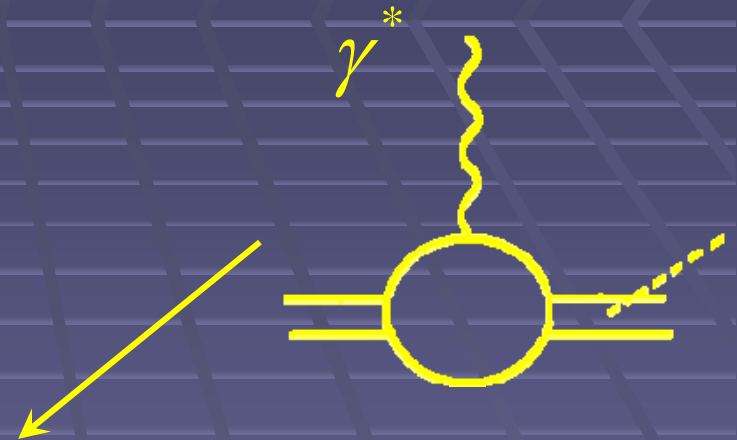
$$F_2(W, Q^2) \text{ at } W - W_{th} - m_\pi$$

$$Q^2 \geq 9 \text{ GeV}^2$$

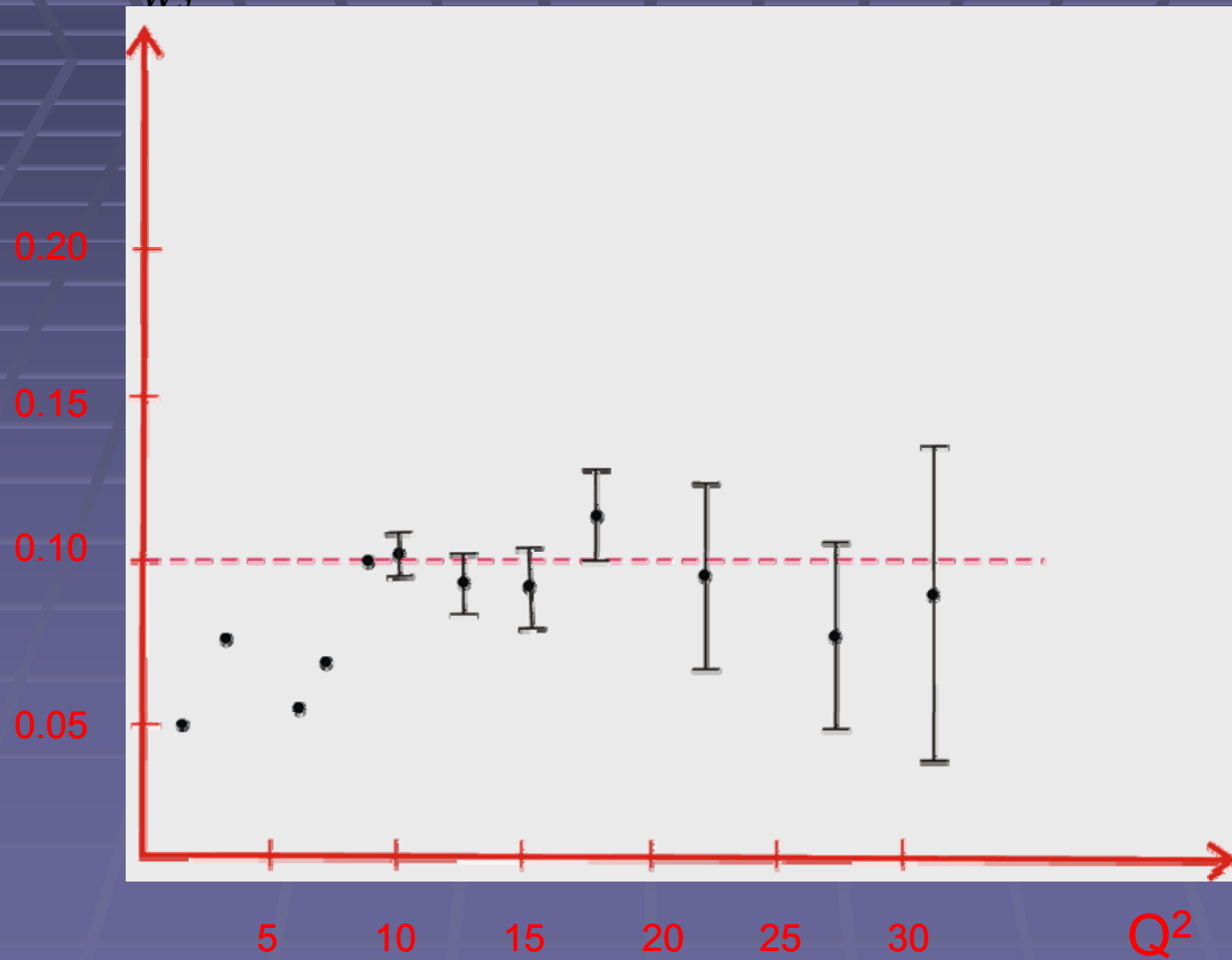
$$F_2^P(W, Q^2) = \frac{Q^2 \beta(w)}{16\pi^2} *$$

$$* \left[ \sum_{\substack{x=p\pi^0 \\ n\pi^+}} \left| A(\gamma^* p \rightarrow x) \right|^2 + \frac{3g_A^2 G_{Mp}^2(Q^2) \beta^2(w) w^4}{4(w^2 - M_N^2 + m_\pi^2)^2} + O\left(\frac{m_\pi}{\Lambda}\right) \right]$$

$$\beta(w) = \sqrt{1 - \frac{(M_N + m_\pi)^2}{W^2}} \sqrt{1 - \frac{(M_N - m_\pi)^2}{W^2}}$$



$$\int_{W_1^2}^{1,4\text{GeV}^2} dw^2 Q^6 F_2^P(W, Q)$$



*For  $Q^2 \geq 9\text{GeV}^2$*

*E136 gives*

$$\int_{W_{th}^2}^{1,4\text{GeV}^2} dw^2 Q^6 F_2(W, Q^2) = [0,1 \pm 0,02] \text{GeV}^8$$

*Nambu – Shrauner theorem*

$$= [0,03 \pm 0,01] \text{GeV}^8$$

*/Nambu, Shrauner PR '62/*

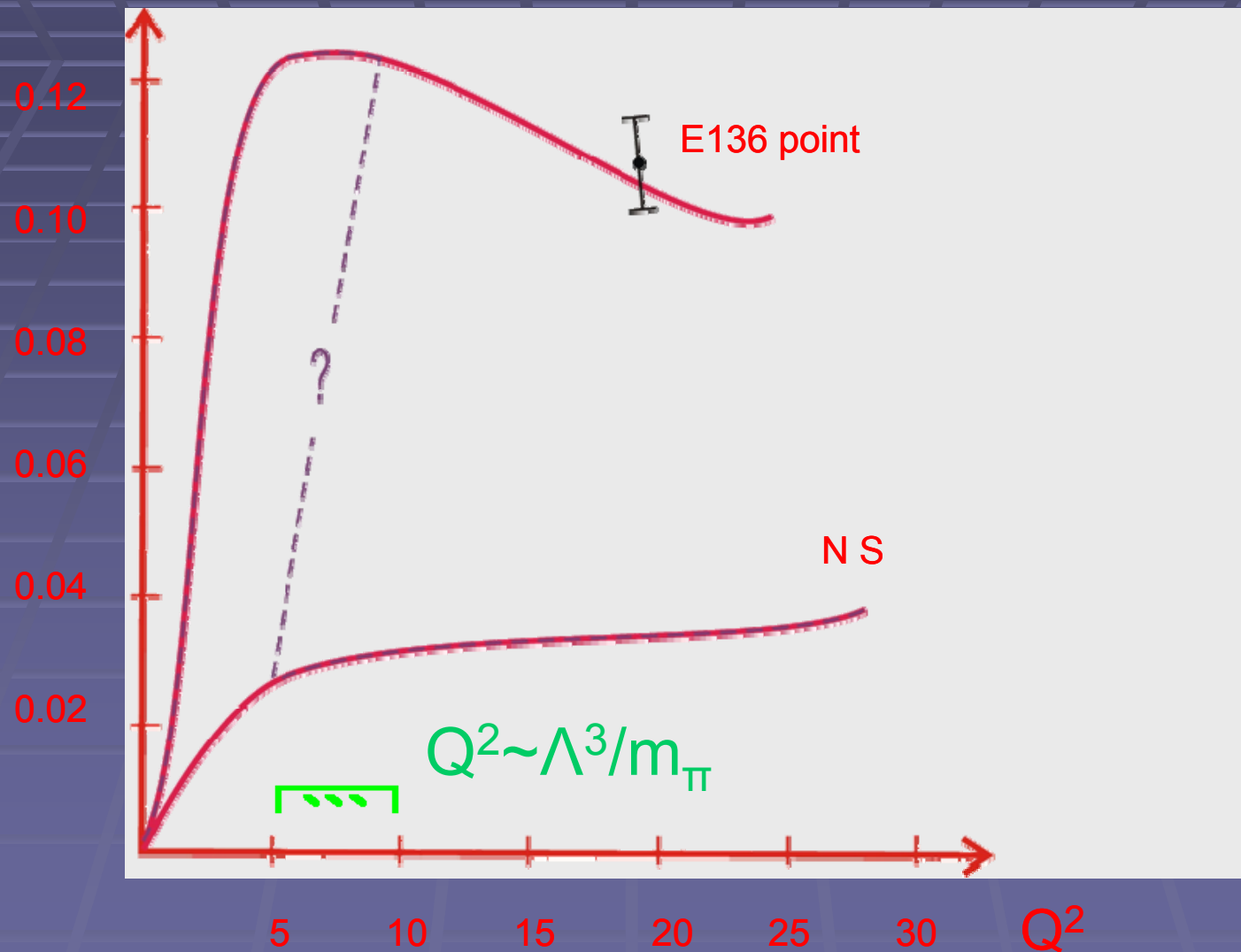
*New soft pion theorem*

$$= [0,11 \pm 0,02] \text{GeV}^8$$

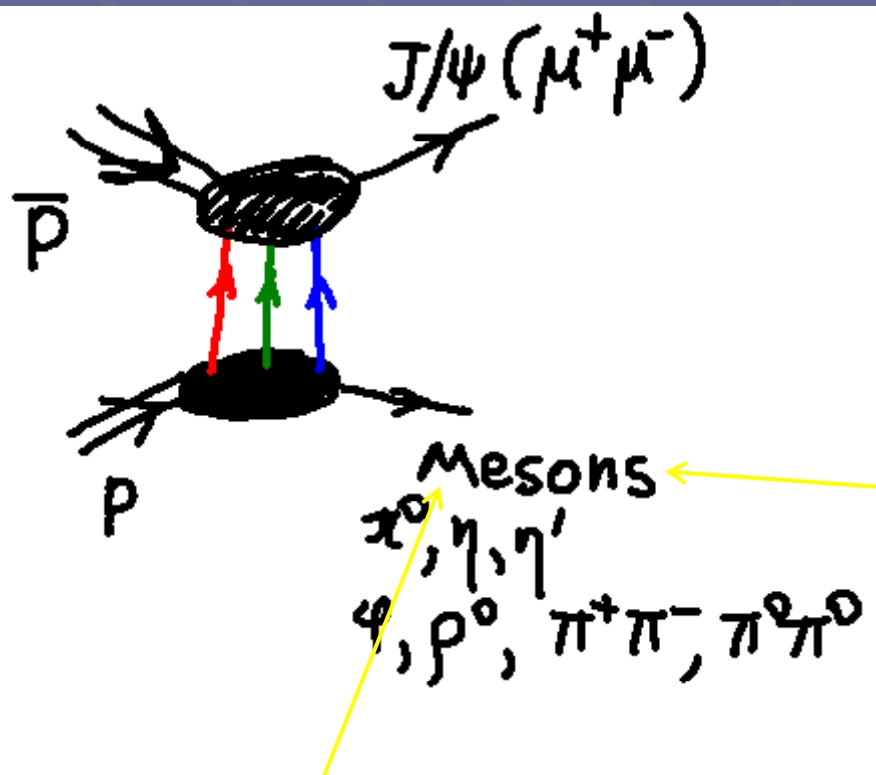
*/Dobuditsa, Strikman, MVP, PRL '01/*

**∴) Nucleon FF's from experiment**

... ? ... very interesting domain Transition  
from one type of SPT<sub>h</sub> to another one



# Opportunities for PANDA



Frankfurt, Strikman, MVP '98

Lansberg, Pire, Szymanowski '07

Soft relative to the proton

Direct measurement of 5+ quark component of nucleon WF

Measuring different mesons we probe different components of nucleon WF

If here is a pion, eta (gluon rich)  $\rightarrow$  chiral dynamics and anomaly for 5 quark cluster

Possibility to figure out how the chiral symmetry is restored

